

## Higher rank Einstein solvmanifolds

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ABSTRACT. In this paper we study the structure of standard Einstein solvmanifolds of arbitrary rank. Also the validity of a variational method for finding standard Einstein solvmanifolds is proved.

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### 1. INTRODUCTION

General form of standard Einstein solvmanifolds were determined by Jense Heber (see [2]). Later, Gorge Lauret deeply studied this kind of manifolds. Solvable Lie group endowed with the left invariant Riemannian metric is called solvmanifold. Let  $S$  be a simply connected Lie group with the corresponding Lie algebra  $s$  endowed with the inner product determined by  $\langle \cdot, \cdot \rangle$  and solvable Lie bracket  $[\cdot, \cdot]$ . We call  $S$  a higher rank solvmanifold if

$$s = n \oplus a; \quad n = [s, s], \quad a = n^\perp,$$

where  $n$  is a metric nilpotent Lie algebra of dimension  $k$ . The codimension  $n$  is called the rank of  $S$ . The solvable Lie group  $(S, [\cdot, \cdot], \langle \cdot, \cdot \rangle)$  is called standard if  $a$  is abelian and it is said to be Einstein if its Ricci tensor  $ric_{[\cdot, \cdot]}$  satisfies  $ric \langle \cdot, \cdot \rangle = c \langle \cdot, \cdot \rangle$ , for some  $c \in \mathbb{R}$ .  $s$  is called a metric solvable extension of  $n$  if the restriction of the Lie bracket and inner product of  $s$  to  $n$  coincide respectively with the Lie bracket and inner product of  $n$ . Let  $\Lambda^2 n^* \otimes n$

be the vector space of all bilinear skew-symmetric maps from  $n \times n$  to  $n$ . There is a natural action of  $GL(k)$  on  $\Lambda^2 n^* \otimes n$  which is given by

$$\phi.\mu(X, Y) = \phi\mu(\phi^{-1}X, \phi^{-1}Y); \quad X, Y \in n, \quad \phi \in GL(k), \quad \mu \in \Lambda^2 n^* \otimes n.$$

Let  $N$  denote a simply connected nilpotent Lie group with Lie algebra  $(n, \mu)$  endowed with the left invariant Riemannian metric  $\langle \cdot, \cdot \rangle_n$ , where  $\mu$  is a nilpotent Lie algebra on  $n$ . The Ricci operator  $R_\mu : n \rightarrow n$  of  $N$  is defined by

$$\langle R_\mu \cdot, \cdot \rangle_n = \text{ric} \langle \cdot, \cdot \rangle_n.$$

This operator is reduced to

$$(1.1) \quad \langle R_\mu X, Y \rangle = -\frac{1}{2} \sum_{i,j} \langle \mu(X, X_i), X_j \rangle \langle \mu(Y, X_i), X_j \rangle \\ + \frac{1}{4} \sum_{i,j} \langle \mu(X_i, X_j), X \rangle \langle \mu(X_i, X_j), Y \rangle,$$

for all  $X, Y \in n$ , where  $\{X_1, X_2, \dots, X_k\}$  is any orthonormal basis of  $n$ . The inner product  $\langle \cdot, \cdot \rangle_n$  determines an inner product on  $\Lambda^2 n^* \otimes n$ , denoted by  $\langle \cdot, \cdot \rangle$  and given by

$$\langle \mu, \lambda \rangle = \sum_{i,j,k} \langle \mu(X_i, X_j), X_k \rangle \langle \lambda(X_i, X_j), X_k \rangle.$$

Also it naturally determines a norm on  $\Lambda^2 n^* \otimes n$  defined by

$$\forall \lambda \in \Lambda^2 n^* \otimes n \quad \|\lambda\| = \sum_{ijk} \langle \lambda(X_i, X_j), X_v \rangle^2.$$

Consider the Riemannian function

$$F_k : \Lambda^2 n^* \otimes n \longrightarrow \mathbb{R}, \quad F_k(\mu) = \text{tr} R_\mu^2$$

and the sphere  $S_r$  given by

$$S_r = \{\mu \in \Lambda^2 n^* \otimes n; \|\mu\|^2 = 2r^2\},$$

for some  $r \in \mathbb{R}$ . Let  $\mathfrak{N}_k$  be the vector space of all nilpotent Lie brackets on  $n$  and  $Der(\mu)$  be the Lie algebra of all derivations on  $n$ . Then  $\mu \in \mathfrak{N}_k$  is called a Ricci soliton if  $R_\mu = cI + D$ , for some  $D \in Der(\mu)$  and  $c \in \mathbb{R}$ .

In [6], Jorge Lauret has proved that the standard Einstein solvmanifolds are exactly the critical points of modified scalar curvature function  $F_k|_{S_1}$ .

**Theorem 1.1.** [6]. *For  $\mu \in \mathfrak{N}_k \cap S_1$ , the following statements are equivalent:*

- (i)  $\mu$  is a critical point of  $F_k|_{S_1}$ .
- (ii)  $\mu$  is a critical point of  $F_k|_{GL(k).\mu \cap S_1}$ .
- (iii)  $\mu$  admits a rank-one extension which is Einstein.
- (iv)  $\mu$  is a Ricci soliton.

2. EINSTEIN SOLVMANIFOLDS OF RANK  $\geq 1$ 

In this section, the structure of standard Einstein solvmanifolds is introduced. The next lemma provides some useful properties of solvable Lie bracket and inner product of a solvable Lie group. Then, we extend the Ricci soliton in [4] for any arbitrary rank. We call it multiple Ricci soliton.

**Lemma 2.1.** [2]. *Let  $(s = n \oplus a, [\cdot, \cdot], \langle \cdot, \cdot \rangle)$  be a metric solvable extension of  $(n, \mu, \langle \cdot, \cdot \rangle)$ , where for every  $0 \neq A \in a$ ,  $ad_A$  is nonzero and symmetric, then*

- (i)  $\langle R_{[\cdot, \cdot]}A, B \rangle = -tr(ad_A ad_B)$ , for all  $A, B \in a$ .
- (ii)  $\langle R_{[\cdot, \cdot]}A, X \rangle = 0$ , for all  $A \in a, X \in n$ .
- (iii)  $R_{[\cdot, \cdot]}|_n = -ad_Z|_n + R_\mu$ , where  $\langle Z, X \rangle = tr(ad_X)$ , for all  $X \in n$ .

**Definition 2.2.**  $0 \neq \mu \in \mathfrak{N}_k$  is called a multiple Ricci soliton of degree  $r$  if

- (a)  $R_\mu = c_\mu I + D_\mu$ ;  $D_\mu \in Der(\mu)$ ,  $c_\mu \in \mathbb{R}$ .
- (b) There are nonzero symmetric derivations  $D_i$ ,  $1 \leq i \leq r$ , such that

$$D_\mu = D_1 + D_2 + \dots + D_r, \quad tr D_i D_j = -\delta_{ij} c_\mu tr D_i.$$

**Remark 2.3.** If  $\mu$  is a multiple Ricci soliton of degree  $r$ , then  $\mu$  is a multiple Ricci soliton of degree less than  $r$ . Therefore,  $\mu$  is the critical point of  $F_k|_{S_r \cap GL(k) \cdot \mu}$ .

Using Definition 2.1, we study the structure of standard Einstein solvmanifolds as follows.

**Proposition 2.4.** *For  $0 \neq \mu \in \mathfrak{N}_k \cap S_r$ , the following statements are equivalent:*

- (i)  $\mu$  admits a metric extension which is Einstein.
- (ii)  $\mu$  is a multiple Ricci soliton.

*Proof.* Let the Lie algebra  $(n, \mu)$  admit an Einstein metric extension  $S$  with corresponding Lie algebra  $(s = n \oplus a, [\cdot, \cdot], \langle \cdot, \cdot \rangle)$  such that  $dim(a) = r$ . Let  $\{H_1, H_2, \dots, H_r\}$  be an orthonormal basis for  $a$  and  $Z$  be the mean curvature vector field for the simply connected Lie group  $N$  with Lie algebra  $n$ . A straightforward calculation shows that  $D_\mu = D_1 + D_2 + \dots + D_r$ , where  $D_\mu := ad_Z|_n$  and  $D_i = tr(ad_{H_i})ad_{H_i}|_n$ .  $[\cdot, \cdot]$  is the Lie bracket. Hence,  $D_i \mu(\cdot, \cdot) = \mu(\cdot, D_i \cdot) + \mu(D_i \cdot, \cdot)$ ; that is to say,  $D_i$ 's are derivations on  $n$ . Suppose that  $D_i$ 's and  $D_\mu$  are symmetric (see [2; 4.10]). Let  $Z_i = tr(ad_{H_i})H_i$ , then Lemma 2.1 implies that

$$tr D_i D_j = tr(ad_{Z_i} ad_{Z_j}) = -\langle R_{[\cdot, \cdot]}Z_i, Z_j \rangle = -c_\mu \langle Z_i, Z_j \rangle = -\delta_{ij} c_\mu tr D_i.$$

Also  $R_\mu = c_\mu I + D_\mu$ , for some  $c_\mu \in \mathbb{R}$ . Therefore,  $\mu$  is a multiple Ricci soliton.

Conversely, let  $\mu$  be a multiple Ricci soliton i.e.

- (a)  $R_\mu = c_\mu I + D_\mu$ ;  $D_\mu \in Der(n, \mu)$ ,  $c_\mu \in \mathbb{R}$ .

(b) There are symmetric derivations  $D_i$ ,  $1 \leq i \leq r$ , such that

$$D_\mu = D_1 + D_2 + \dots + D_r, \quad \text{tr} D_i D_j = -\delta_{ij} c_\mu \text{tr} D_i.$$

Let  $(n, \mu, \langle \cdot, \cdot \rangle_n)$  be a Lie algebra with orthonormal basis  $\{X_1, X_2, \dots, X_k\}$ . We define Lie algebra  $s$  with a simply connected Lie group  $S$  as follows

$$s = n \oplus \sum_i \mathbb{R} Z_i,$$

endowed with the inner product  $\langle \cdot, \cdot \rangle$  defined by

$$\langle Z_i, Z_j \rangle = \delta_{ij} \text{tr} D_i, \quad \langle Z_i, n \rangle = 0, \quad \langle \cdot, \cdot \rangle|_{n \times n} = \langle \cdot, \cdot \rangle_n.$$

Also, Lie bracket  $[\cdot, \cdot]$  is defined by

$$[Z_i, Z_j] = 0, \quad [Z_i, X_j] = -[X_j, Z_i] = D_i X_j, \quad [\cdot, \cdot]|_{n \times n} = \mu.$$

Clearly  $[\cdot, \cdot]$  is a Lie bracket, since  $D_i$ 's are derivations.  $\{D_\mu X_1, D_\mu X_2, \dots, D_\mu X_k\}$  is a linearly independent set which generates a subalgebra of  $[s, s]$ . Therefore,  $n = [s, s]$ .  $\mu$  is nilpotent hence  $[\cdot, \cdot]$  is a solvable Lie bracket. Finally using Lemma 2.1, we have

$$\begin{aligned} \langle R_{[\cdot, \cdot]} Z_i, Z_j \rangle &= -\text{tr}(D_i D_j) = \delta_{ij} c_\mu \text{tr} D_i = c_\mu \langle Z_i, Z_j \rangle, \quad \langle R_{[\cdot, \cdot]} Z_i, n \rangle = 0, \\ \langle R_{[\cdot, \cdot]} X_i, X_j \rangle &= \langle (-D_\mu + R_\mu) X_i, X_j \rangle = \langle c_\mu X_i, X_j \rangle = c_\mu \langle X_i, X_j \rangle, \end{aligned}$$

which implies that  $\langle \cdot, \cdot \rangle_s$  is a Einstein metric. This completes the proof.  $\square$

Using Proposition 2.1, we get a higher rank Einstein solvmanifold as the direct sum of the Lie algebras.

**Proposition 2.5.** *If nonzero nilpotent Lie brackets  $\mu_1$  and  $\mu_2$  are Ricci solitons, then  $\mu = \mu_1 \oplus \mu_2$  is a multiple Ricci soliton of degree 2.*

*Proof.*  $\mu_1$  and  $\mu_2$  are Ricci solitons i.e.

$$(2.1) \quad R_{\mu_i} = c_{\mu_i} I + D_{\mu_i}; \quad D_{\mu_i} \in \text{Der}(\mu_i), \quad c_{\mu_i} \in \mathbb{R}, \quad i = 1, 2.$$

Up to isometry and scaling we can determine norms of  $\mu_1$  and  $\mu_2$  such that  $c_{\mu_1} = c_{\mu_2}$ . Set

$$D_\mu = \begin{bmatrix} D_{\mu_1} & 0 \\ 0 & D_{\mu_2} \end{bmatrix},$$

$$R_\mu = \begin{bmatrix} R_{\mu_1} & 0 \\ 0 & R_{\mu_2} \end{bmatrix},$$

Then  $R_\mu = c_{\mu_1} I + D_\mu$ ,  $R_\mu = R_{\mu_1} \oplus R_{\mu_2}$  and  $D_\mu = D_{\mu_1} \oplus D_{\mu_2}$ . Also, by Theorem 1.3,  $\text{tr} D_{\mu_i} D_{\mu_j} = -\delta_{ij} c_{\mu_1} \text{tr} D_{\mu_i}$ ;  $i = 1, 2$ . Therefore  $\mu$  is a multiple Ricci soliton which admits a 2-rank Einstein solvable extension.  $\square$

**Corollary 2.6.** *If nonzero nilpotent Lie brackets  $\mu_i$ 's,  $1 \leq i \leq r$ , are Ricci solitons, then  $\mu = \mu_1 \oplus \mu_2 \oplus \dots \oplus \mu_r$  is a multiple Ricci soliton of degree  $r$  which admits an Einstein solvable extension of rank  $r$ .*

**Remark 2.7.** There exist 31 Ricci soliton nonzero Lie algebras of dimension 6 (see [7]), which by direct sum of them, we can obtain a lot of multiple Ricci soliton nilpotent Lie algebras.

### 3. STANDARD METHODS

The goal of this section is to present certain results from [4] and [6]. In view of [4], Jorge Lauret has used a variational method for finding standard Einstein solvmanifolds. We will demonstrate this method in Theorem 3.1. We first give some preliminaries.

**Lemma 3.1.** (*Lagrange multiplier theorem*) [1]. *Let  $P$  and  $M$  be smooth manifolds and  $g : M \rightarrow P$  be a smooth submersion. Let  $f : M \rightarrow \mathbb{R}$  be  $C^r$ ,  $m \in M$  and  $p \in P$  such that  $m \in g^{-1}(p)$ , then the following statements are equivalent:*

- (i)  $m$  is a critical point of  $f|_{g^{-1}(p)}$ .
- (ii) There are  $\lambda \in T_p^*P$  such that  $T_m f = \lambda \circ T_m g$ .

The vector space  $\mathfrak{N}_k$  is  $GL(k)$ -invariant, so we can refine Theorem 1.1 more accurately as follows.

**Lemma 3.2.** *For  $0 \neq \mu \in \mathfrak{N}_k$  and  $\psi \in GL(k)$  the following statements are equivalent:*

- (i)  $\psi \cdot \mu$  is a Ricci soliton.
- (ii)  $\psi \cdot \mu$  is a critical point of  $F_k|_{S_r \cap GL(k) \cdot \mu}$ .
- (iii)  $\psi$  is a solution of the following system of equations:

$$\begin{cases} \|\phi \cdot \mu\|^2 = 2r^2 \\ \frac{\partial F_k(\phi \cdot \mu)}{\partial \phi_{ij}} = t \frac{\partial (\|\phi \cdot \mu\|)}{\partial \phi_{ij}} \end{cases}$$

where  $t \in \mathbb{R}$  and  $\phi \in GL(k)$ .

*Proof.* Let  $g(\lambda) = \frac{1}{2}\|\lambda\|^2$  be a function on  $\Lambda^2 n^* \otimes n$  and use Theorem 1.1 and Lemma 3.1.

It is possible that the above system of equations is not solvable, hence we assume that for some  $k \in \mathbb{N}$  and every  $\phi \in GL(k)$  there exists  $\mu \in \mathfrak{N}_k$  such that  $\phi \cdot \mu$  isn't a Ricci soliton.  $\square$

**Notation 3.3.** *Suppose that  $DGL(k) := \{\phi \in GL(k); \phi \text{ is diagonal}\}$ , For any  $\mu \in \mathfrak{N}_k \subseteq \Lambda^2 n^* \otimes n$ , set*

$$\mu(X_i, X_j) = \sum_v c_{ijv} X_v, \quad \phi = \text{diag}(\phi_1, \phi_2, \dots, \phi_k),$$

then

$$(\phi \cdot \mu)(X_i, X_j) = \sum_{x_{ijv} \neq 0} x_{ijv} X_v; \quad x_{ijv} = \frac{\phi_v c_{ijv}}{\phi_i \phi_j}$$

and for any  $i, j$  and  $v$  such that  $\langle \mu(X_i, X_j), X_v \rangle \neq 0$ , the diagonal elements of  $R_{\phi, \mu}$  are equal to

$$\begin{aligned} (R_{\phi, \mu})_{ii} &= -\frac{1}{2}(-x_{ijv}^2 + \sum_{\substack{rst \\ x_{rst} \neq x_{ijv}}} \delta^{rst, i} x_{rst}^2), \\ (R_{\phi, \mu})_{jj} &= -\frac{1}{2}(-x_{ijv}^2 + \sum_{\substack{rst \\ x_{rst} \neq x_{ijv}}} \delta^{rst, j} x_{rst}^2), \\ (R_{\phi, \mu})_{vv} &= \frac{1}{2}(x_{ijv}^2 + \sum_{\substack{rst \\ x_{rst} \neq x_{ijv}}} \delta^{rst, v} x_{rst}^2), \end{aligned}$$

where  $\delta^{rst, i}$ ,  $\delta^{rst, j}$  and  $\delta^{rst, v}$  are equal to 0, 1 or -1.

**Lemma 3.4.** [6]. Let  $(n, \mu, \langle \cdot, \cdot \rangle)$  be a Lie algebra and  $P_\mu = \text{Sym}(n) \cap \text{Der}(n)$ , then  $R_\mu \perp P_\mu$  with inner product  $\text{tr}(AB)$  on  $\text{Sym}(n) \times \text{Sym}(n)$ .

**Theorem 3.5.** (Lauret theory) For every  $\phi \in GL(k)$  and  $\mu \in \mathfrak{N}_k$  if  $\phi, \mu \in S_r$  and  $R_\mu$  is diagonal, then for any  $i, j$  and  $v$  such that  $\langle \mu(X_i, X_j), X_v \rangle \neq 0$ , the following statements are equivalent:

- (i)  $R_{\psi, \mu} = c_{\psi, \mu} I + D_{\psi, \mu}$ ;  $D_{\psi, \mu} \in \text{Der}(\psi, \mu)$
- (ii)  $c_{\psi, \mu}$  and  $a_{ijv}^2$ 's are solutions of the system

$$\begin{cases} \sum_{i, j, v} x_{ijv}^2 = r^2 \\ \frac{\partial F_k(\phi, \mu)}{\partial u_{ijv}} \Big|_{x_{ijv} := a_{ijv}} = -c_{\psi, \mu} \end{cases}$$

where  $(\phi, \mu)(X_i, X_j) = \sum_{x_{ijv} \neq 0} x_{ijv} X_v$ ,  $u_{ijv} = x_{ijv}^2$ ,  $\psi, \mu = \phi, \mu|_{x_{ijv} := a_{ijv}}$ .

*Proof.* By Lemma 3.2 and the chain rule, it is easy to see that  $\psi, \mu$  is a Ricci soliton if and only if variation  $t$  and  $a_{ijv}^2$ 's are solutions of the system

$$\begin{cases} \sum_{i, j, v} x_{ijv}^2 = r^2 \\ \frac{\partial F_k(\phi, \mu)}{\partial u_{ijv}} \Big|_{x_{ijv} := a_{ijv}} = t \end{cases}$$

Now we shall obtain the Lagrangian coefficient. By Lemma 3.3 it is easy to see that  $\text{tr} R_\mu^2 = c_\mu \text{tr} R_\mu$ . Also  $\text{tr} R_\mu = -\frac{1}{2} \|\mu\|^2$ . Thus  $F_k = -c_\mu r^2$ . Consequently  $\frac{\partial F_k(\phi, \mu)}{\partial u_{ijv}} \Big|_{x_{ijv} := a_{ijv}} = -c_\mu$ .  $\square$

Finally, we exhibit a rank-two Einstein solvmanifold of dimension 8 and a rank-three Einstein solvmanifold of dimension 15.

**Example 3.6.** Let  $\mu = \mu_1 \oplus \mu_2$ , where  $\mu_1(X_1, X_2) = X_5$ ,  $\mu_2(X_3, X_4) = X_6$  and  $\phi = \text{diag}(\phi_1, \phi_2, \dots, \phi_6)$ , then

$$\phi, \mu_1(X_1, X_2) = \frac{\phi_5}{\phi_1 \phi_2} X_5 \quad \text{and} \quad \phi, \mu_2(X_3, X_4) = \frac{\phi_6}{\phi_3 \phi_4} X_6.$$

Let  $x := \frac{\phi_5}{\phi_1\phi_2}$ ,  $y := \frac{\phi_6}{\phi_3\phi_4}$ . Using Theorem 3.1, it is easy to see  $x^2 = y^2 = 1$ . Therefore  $D_{\mu_1} = \text{diag}(1, 1, 0, 0, 2, 0)$  and  $D_{\mu_2} = \text{diag}(0, 0, 1, 1, 0, 2)$ . If  $(n_1, \mu_1, \langle \cdot, \cdot \rangle_1)$  and  $(n_2, \mu_2, \langle \cdot, \cdot \rangle_2)$  are nilradical Lie algebras, define the Lie algebra  $s$  with simply connected Lie group  $S$  using the following direct sum

$$s = \mathbb{R}Z_1 \oplus n_1 \oplus \mathbb{R}Z_2 \oplus n_2$$

endowed with the inner product  $\langle \cdot, \cdot \rangle$  defined by

$$\langle Z_i, Z_j \rangle = 4\delta_{ij}, \langle Z_i, n \rangle = 0, \langle \cdot, \cdot \rangle|_{n_i \times n_i} = \langle \cdot, \cdot \rangle_i; \quad 1 \leq i, j \leq 2,$$

where the Lie bracket  $[\cdot, \cdot]$  on  $s$  is defined by

$$\begin{aligned} [Z_1, X_1] &= -[X_1, Z_1] = X_1, & [Z_1, X_2] &= -[X_2, Z_1] = X_2, \\ [Z_1, X_5] &= -[X_5, Z_1] = 2X_5, & [Z_2, X_3] &= -[X_3, Z_2] = X_3, \\ [Z_2, X_4] &= -[X_4, Z_2] = X_4, & [Z_2, X_6] &= -[X_6, Z_2] = 2X_6, \\ [Z_i, Z_j] &= 0, & [\cdot, \cdot]_s|_{n_i \times n_i} &= \mu_i; \quad 1 \leq i, j \leq 2 \end{aligned}$$

and it is equal to zero otherwise. It is easy to see that  $R_{[\cdot, \cdot]_s} = -\frac{3}{2}I_{8 \times 8}$  which implies that  $S_\mu$  is Einstein of rank 2.

**Example 3.7.** Let  $\mu = \mu_1 \oplus \mu_2 \oplus \mu_3$ , where  $\mu_1$  and  $\mu_2$  are Lie algebras given in Example 1 and  $\mu_3 \in \mathfrak{N}_6$  is given by

$$\mu_3(X_7, X_i) = X_{i+1}; \quad 8 \leq i \leq 11.$$

Every  $\phi.\mu_3 \in DGL(6).\mu_3 \cap S_{\sqrt{\frac{30}{13}}}$  is equal to

$$\phi.\mu_3(X_7, X_i) = a_{7,i,i+1}X_{i+1}; \quad i = 8, 9, 10, 11.$$

By Lauret theory it is easy to see that a critical point of  $F_6$  restricted to the leaf  $\sum_{7 < i < 12} a_{7,i,i+1}^2 = \frac{20}{13}$  is equal to  $\mu_3\{a_{7,i,i+1}\}$  where

$$a_{7,8,9}^2 = \frac{12}{13}, \quad a_{7,9,10}^2 = \frac{9}{13}, \quad a_{7,10,11}^2 = \frac{3}{13}, \quad a_{7,11,12}^2 = \frac{9}{13}, \quad c_{\psi.\mu_3} = -\frac{3}{2}$$

and

$$D_{\mu_3(a_{7,i,i+1})} = \text{diag}(0, 0, 0, 0, 0, 0, \frac{15}{26}, \frac{42}{26}, \frac{36}{26}, \frac{51}{26}, \frac{48}{26}, \frac{51}{26}).$$

Let  $(n_3, \mu_3(a_{7,i,i+1}), \langle \cdot, \cdot \rangle_3)$  be a Ricci soliton. Define Lie algebra  $s$  with simply connected Lie group  $S$  using the following direct sum

$$s = \mathbb{R}Z_1 \oplus n_1 \oplus \mathbb{R}Z_2 \oplus n_2 \oplus \mathbb{R}Z_3 \oplus n_3$$

endowed with the inner product  $\langle \cdot, \cdot \rangle$  which is defined by

$$\langle Z_1, Z_1 \rangle = 4, \langle Z_2, Z_2 \rangle = 4, \langle Z_3, Z_3 \rangle = \frac{243}{13}, \langle \cdot, \cdot \rangle|_{n_i \times n_i} = \langle \cdot, \cdot \rangle_i$$

and it is equal to zero otherwise. Lie bracket  $[\cdot, \cdot]$  on  $S$  defined by

$$\begin{aligned} [Z_1, X_1] &= -[X_1, Z_1] = X_1, & [Z_1, X_2] &= -[X_2, Z_1] = X_2, \\ [Z_1, X_5] &= -[X_5, Z_1] = 2X_5, & [Z_2, X_3] &= -[X_3, Z_2] = X_3, \\ [Z_2, X_4] &= -[X_4, Z_2] = X_4, & [Z_2, X_6] &= -[X_6, Z_2] = 2X_6, \end{aligned}$$

$$\begin{aligned}
[Z_3, X_7] &= -[X_7, Z_3] = \frac{15}{26}X_7, & [Z_3, X_8] &= -[X_8, Z_3] = \frac{42}{26}X_8, \\
[Z_3, X_9] &= -[X_9, Z_3] = \frac{36}{26}X_9, & [Z_3, X_{10}] &= -[X_{10}, Z_3] = \frac{51}{26}X_{10}, \\
[Z_3, X_{11}] &= -[X_{11}, Z_3] = \frac{48}{26}X_{11}, & [Z_3, X_{12}] &= -[X_{12}, Z_3] = \frac{51}{26}X_{12}, \\
[\cdot, \cdot]_{n_i \times n_i} &= \mu_{n_i}; \quad i = 1, 2
\end{aligned}$$

and otherwise is equal to zero. It is easy to check that  $R_{[\cdot, \cdot]} = -\frac{3}{2}I_{15 \times 15}$  which implies that  $S_\mu$  is Einstein of rank 3.

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